



Applications of fiber bundles in physics: From cosmology to electronic states of matter

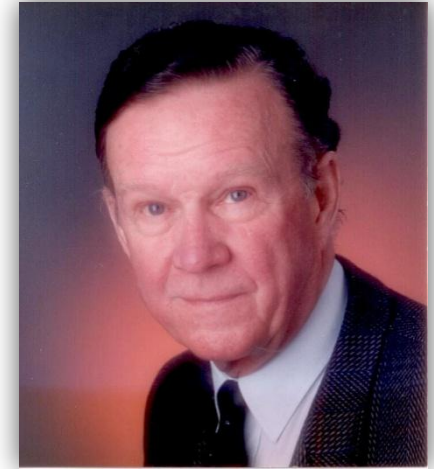
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Colloquium in honor of Prof. Dr. Erik W. Grafarend on occasion of his 75th birthday

Berlin 13 February 2015

Prolog



Erik W. Grafarend and H-RT:

Ekkehart Kröner (1919 – 2000):

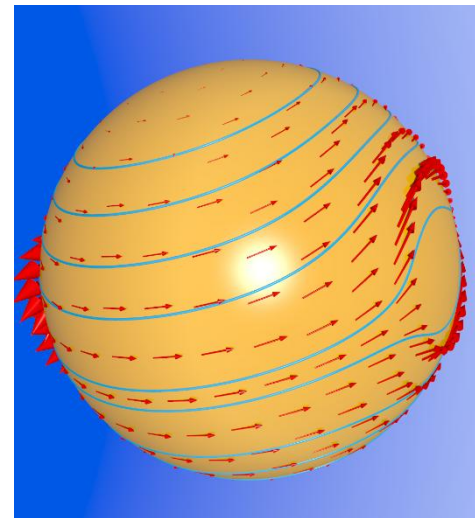
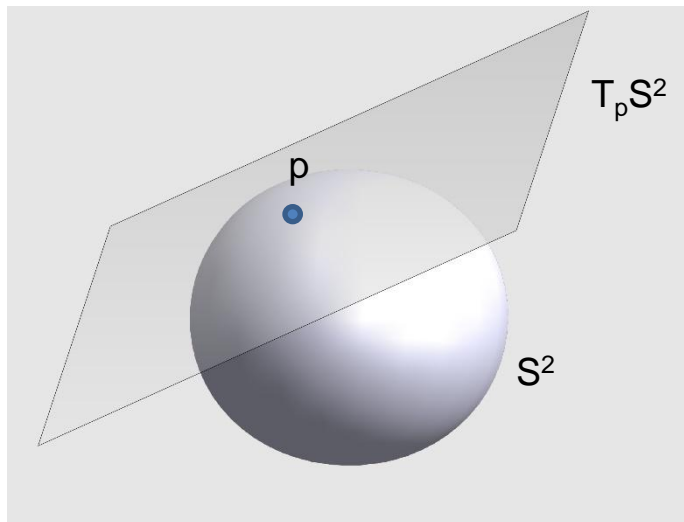
- Lehrstuhl für Theoretische Physik, Clausthal (1963 – 1969)
- Lehrstuhl für Theoretische und Angewandte Physik, Stuttgart (1969 – 1985)

SFB 228 Hochgenaue Navigation – Integration navigatorischer und geodätischer Methoden

The tangent bundle of the sphere

Sphere S^2 : Prominent example of a differential manifold

- Sphere plus tangent planes: tangent bundle TS^2
- TS^2 is example for a fiber bundle:
Basis manifold M is S^2 , fibers are the tangent planes $T_p S^2$
- Vector field: section of the fiber bundle



Parallel transport of vectors on S^2

- Comparison of different tangent spaces by parallel transport along a path

- Levi Civita connection

- Covariant derivative:

$$D_\lambda w := e_\mu (\partial_\lambda w^\mu + \Gamma_{\nu\lambda}^\mu w^\nu)$$

$\Gamma_{\nu\lambda}^\mu$: Connection coefficients

Parallel transport: $D_\lambda w = 0$

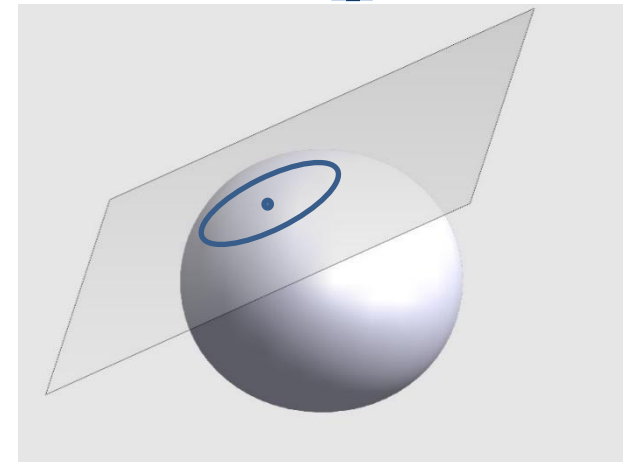
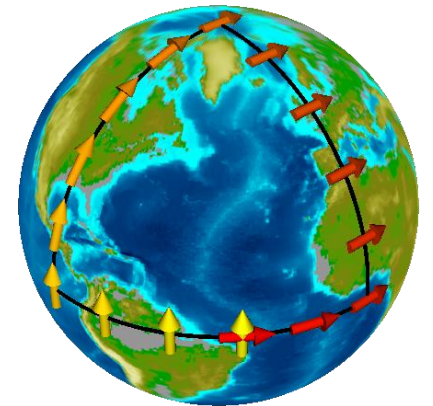
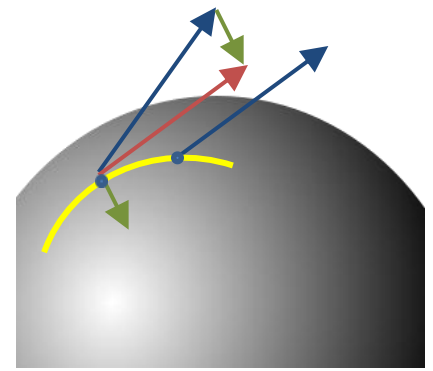
- Curvature: transport along closed path rotates vector (**holonomy**)

- Rotation angle:

$$\Delta\omega = \int_{\text{encircled area}} d\Omega \kappa(\theta, \varphi)$$

κ = Gaussian curvature

- For description of parallel transport and curvature: attach rotation group $SO(2)$ at each point, yields **principal bundle** $(S^2, SO(2)) = (S^2, S^1)$



Fiber bundles

- Consist of a basic differential manifold M
- At each point attached: fiber F_p which is either copy of a vector space (“vector bundle”) or of a (gauge) group (“principle bundle”)
- Prescriptions for glueing the fibers together e.g. Moebius strip (S^1, \mathbb{R})
- Prescription for parallel transport : covariant derivative and curvature

$$D_\lambda w := e_j (\partial_\lambda w^j + \Gamma_{i\lambda}^j w^i) = (\partial_\lambda + \Gamma_{\lambda}) w$$

Γ_{λ} : Connection coefficient matrix

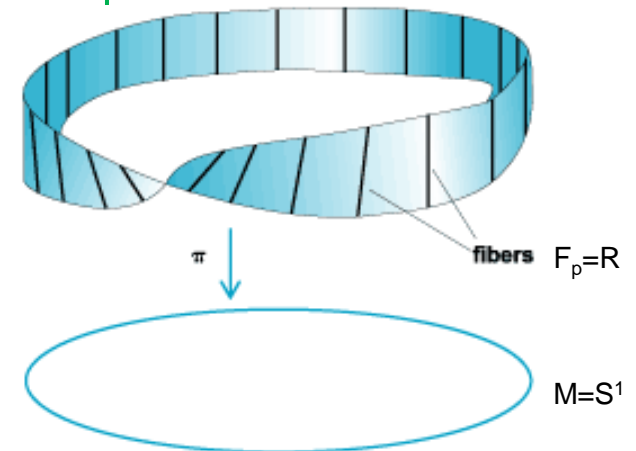
Curvature tensor :

$$R_{j\mu\nu}^i = \partial_\mu \Gamma_{j\nu}^i - \partial_\nu \Gamma_{j\mu}^i + \Gamma_{l\mu}^i \Gamma_{j\nu}^l - \Gamma_{l\nu}^i \Gamma_{j\mu}^l$$

$$R_{\mu\nu} = \partial_\mu \Gamma_{\nu} - \partial_\nu \Gamma_{\mu} + [\Gamma_{\mu}, \Gamma_{\nu}]$$

- Topological quantum numbers, e.g. for 2d tangent bundles: Euler characteristic:

$$\chi = \frac{1}{2\pi} \int_{M^2} d\Omega \kappa(\theta, \varphi) = 2 - g$$



Fiber bundle in cosmology

$$T(\mathbb{R} \oplus \mathbb{R}^3) \text{ or } (\mathbb{R} \oplus \mathbb{R}^3, \text{SO}(1,3))$$



Fiber bundle in electrodynamics

Classical relativistic physics:

- Four momentum: $p^\mu = \begin{bmatrix} E/c \\ \vec{p} \end{bmatrix}$, four velocity: $v^\mu = p^\mu/m$
- In electromagnetic field: $v^\mu = \frac{1}{m}(p^\mu - qA^\mu)$, $A^\mu = \begin{bmatrix} \Phi \\ \vec{A} \end{bmatrix}$

Quantum mechanics:

- $\psi \in \mathbb{C}$, i.e. section through a $(\mathbb{R} \oplus \mathbb{R}^3, \mathbb{C})$ -bundle
- Parallelism?

$$p_\mu \rightarrow \frac{\hbar}{i} \partial_\mu, \quad v_\mu = \frac{\hbar}{im} \left(\partial_\mu - i \frac{q}{\hbar} A_\mu \right) \equiv \frac{\hbar}{im} D_\mu$$

$$D_\mu \psi = 0 \Rightarrow \psi = \psi_0 \exp \left\{ \underbrace{i \frac{q}{\hbar} \int dx^\nu A_\nu}_{\text{phase change} \in U(1)} \right\}$$

- $(\mathbb{R} \oplus \mathbb{R}^3, U(1))$ principle bundle

Topology:

- Curvature equals the electromagnetic field tensor

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu \text{ cp. } R_{\mu\nu} = \partial_\mu \Gamma_{\nu} - \partial_\nu \Gamma_{\mu} + [\Gamma_{\mu}, \Gamma_{\nu}]$$

$$F_{\mu\nu} = \begin{bmatrix} 0 & -E_x & -E_y & -E_z \\ E_x & 0 & B_z & -B_y \\ E_y & -B_z & 0 & B_x \\ E_z & B_y & -B_x & 0 \end{bmatrix}$$

- Topological quantum number: First Chern number

$$\text{ch}^{(1)} = \frac{1}{2\pi} \underbrace{\int dx^1 dx^2 F_{12}}_{\text{2d closed submanifold}}$$

Magnetic monopole of strength γ

- Vector potential and electromagnetic field tensor (acting on C):

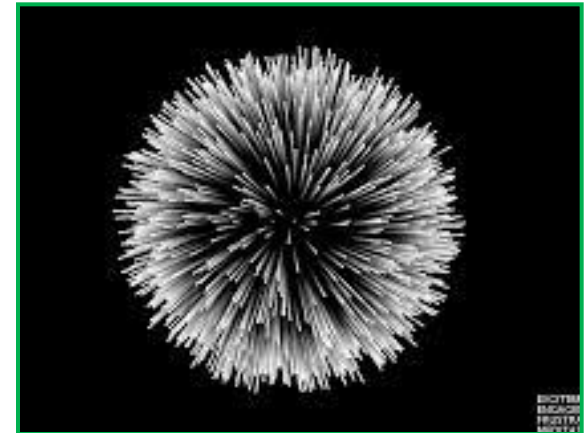
$$-iA_\theta = 0, \quad -iA_\varphi = \gamma i \cos\theta, \quad iF_{\theta\varphi} = \gamma i \sin\theta$$

- Cp. Levi-Civita connection on the sphere in orthonormal basis (acting on R^2):

$$\hat{\Gamma}_{\theta\theta} = 0, \quad \hat{\Gamma}_{\varphi\varphi} = \underbrace{\begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}}_{i} \cos\theta, \quad \hat{R}_{\theta\varphi} = \underbrace{\begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}}_{i} \sin\theta$$

- First chern number: $ch^{(1)} = \frac{1}{2\pi} \int_0^{2\pi} d\varphi \int_0^\pi d\theta F_{\theta\varphi} = 2\gamma$
(R)

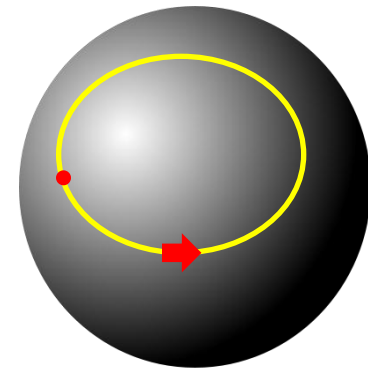
- $(R \oplus R^3, U(1))$: Electromagnetism
- $(R \oplus R^3, U(1) \otimes SU(2))$: Electroweak interaction
- $(R \oplus R^3, SU(3))$: Strong interaction



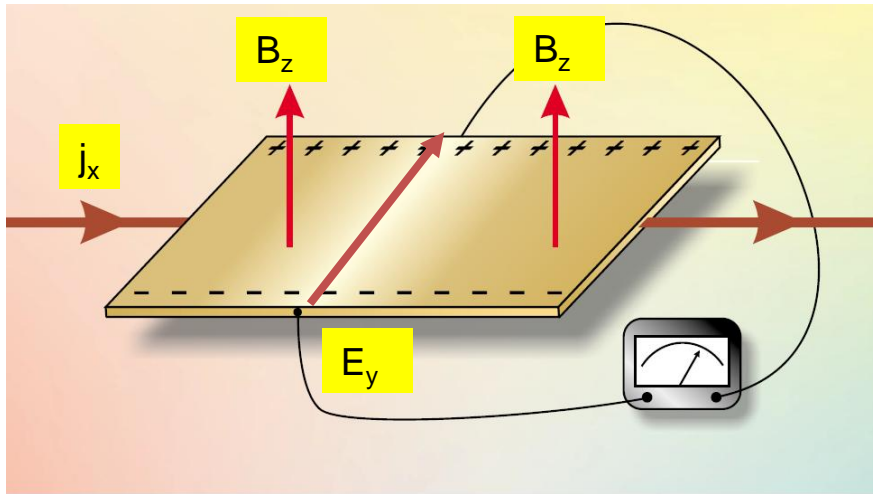
The Berry phase

- Quantum mechanical ground states dependent on parameter $\xi : |m(\xi)\rangle$
- $\xi \in M$ parameter manifold
- Example: Spin-1/2 – particle in magnetic field of fixed strength, but arbitrary orientation : $M = S^2$, $\xi = (\vartheta, \varphi)$
- $|m(\xi)\rangle$ is determined up to a phase factor $e^{i\alpha}$, system is fiber bundle $(M, U(1))$

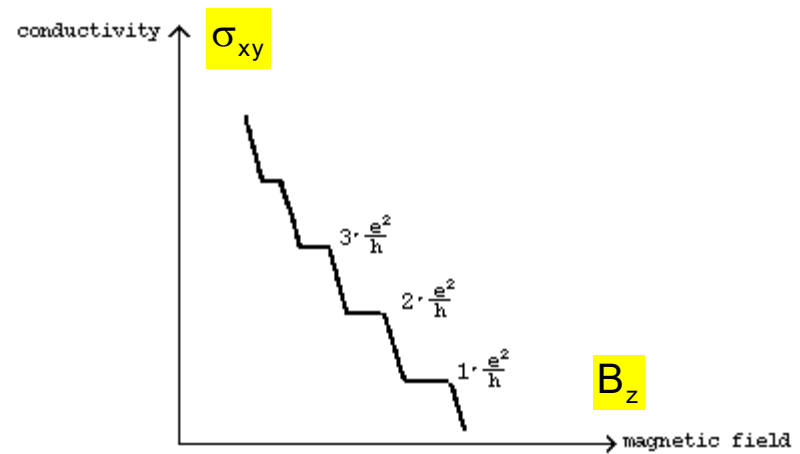
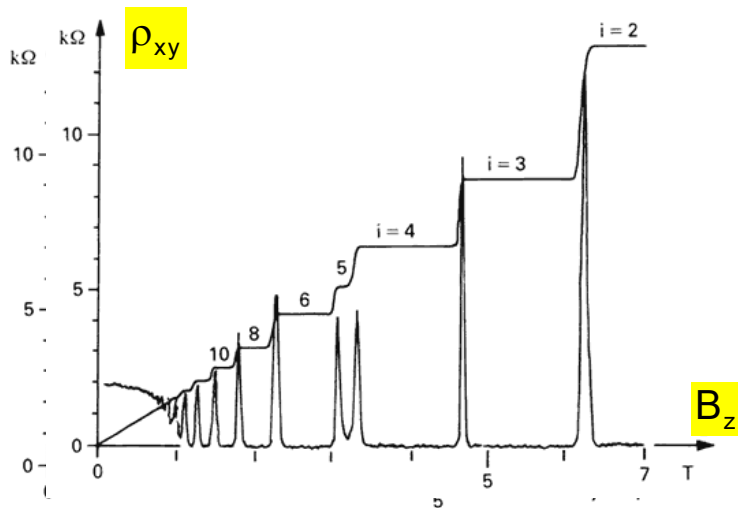
- Adiabatic motion and energy gap :
particles remain in instantaneous state $|m(\xi)\rangle$
- Circular motion: phase change, holonomy
- Covariant derivative $D_\mu = \partial_\mu - iA_\mu$ with
Berry connection $A_\mu = i \langle m | \partial_\mu m \rangle$
Berry curvature $F_{\mu\nu} = i \{ \langle \partial_\mu m | \partial_\nu m \rangle - \langle \partial_\nu m | \partial_\mu m \rangle \}$
- First Chern number for 2d parameter space
$$ch^{(1)}(M) = \frac{1}{2\pi} \int_M dx^1 dx^2 F_{12}$$
- Example spin-1/2 - particle :
 $-iA_\vartheta = 0$, $-iA_\varphi = i \frac{1}{2} \cos \theta$, $iF_{\vartheta\varphi} = i \frac{1}{2} \sin \theta$, $ch^{(1)} = 1$



The Quantum Hall Effect (von Klitzing 1980)

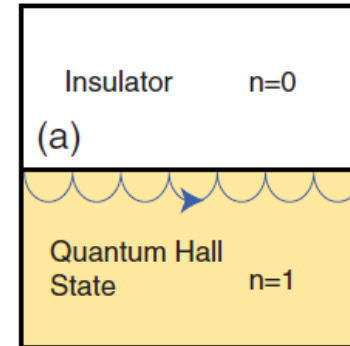


- 2d crystal in magnetic field
- $E_y = \rho_{yx} j_x$ $j_x = \sigma_{xy} E_y$
- $\rho_{yx} = \frac{1}{n} \frac{h}{e^2}$ $\sigma_{xy} = \frac{e^2}{h} n$

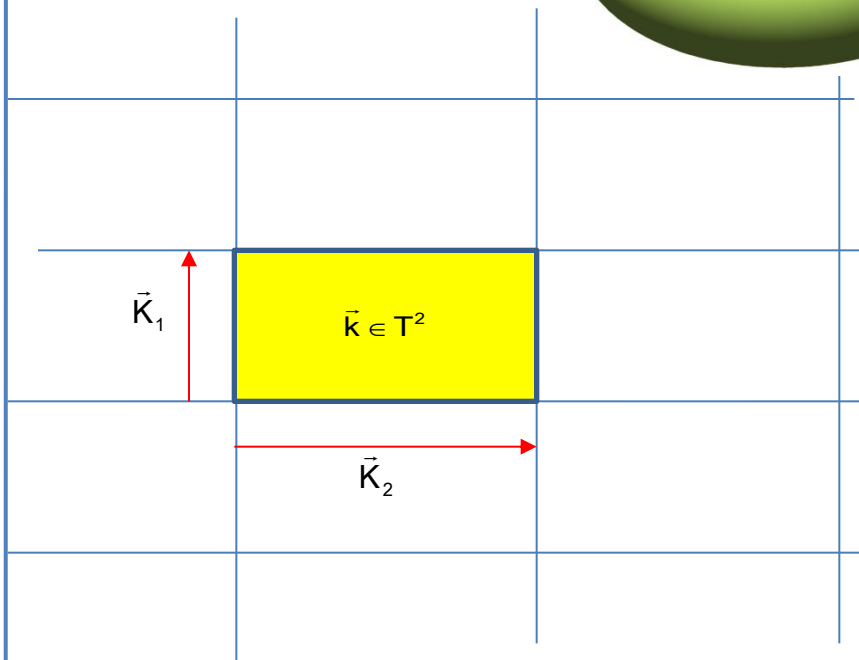


The Quantum Hall Effect

- Wave functions labelled by wavevector $\vec{k} = (k_1, k_2)$: $|u(\vec{k})\rangle$
- Wavevector space is also periodic: $\vec{k} \in \mathbb{R}^2 \text{ mod } (\vec{K}_1, \vec{K}_2) = \mathbb{T}^2$
- Principle bundle $(\mathbb{T}^2, U(1))$



Hasan MZ, Kane CL 2010 RMP 82, 3045



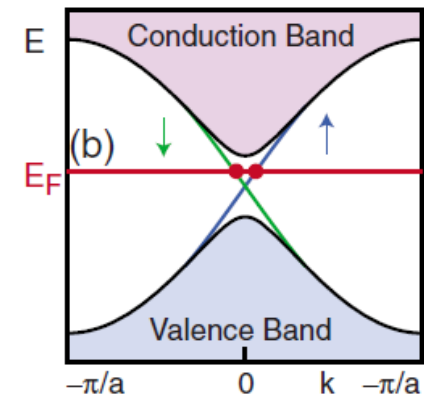
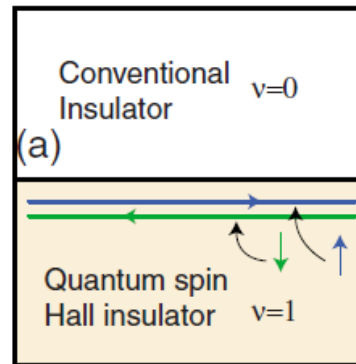
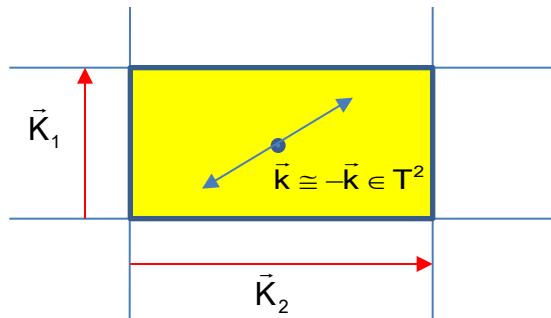
- Kubo transport formula for Hall conductivity :

$$\sigma_{xy} = \frac{e^2}{h} \sum_{\text{occupied bands}} \frac{1}{2\pi} \int_{\text{Torus}} d^2k \underbrace{i\{\langle \partial_1 u | \partial_2 u \rangle - \langle \partial_2 u | \partial_1 u \rangle\}}_{\text{Berry curvature } F_{12}}$$

$$\sigma_{xy} = \frac{e^2}{h} \underbrace{\sum_{\text{occupied bands}} \text{ch}^{(1)}(\mathbb{T}^2, U(1))}_n$$

The spin quantum Hall effect

- Restriction on the base manifold by time reversal symmetry: $\vec{k} \cong -\vec{k}$
- New base manifold T^2 / \cong
- New classification by $\nu \in \mathbb{Z}_2 = \{0,1\}$
- HgTe/CdTe quantum well structures in 2d
- Spin polarized edge currents
- $\text{Bi}_{1-x}\text{Sb}_x$ in 3d, surface currents



Hasan MZ, Kane CL 2010 RMP 82, 3045

Summary

- Up to 1980: Quantum numbers based on symmetry
- Easy to break, lift of degeneracies
- Since 1980: Topological quantum numbers
- Robust

Congratulations and best wishes, Erik



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